

NPS-57Fu72051A

NAVAL POSTGRADUATE SCHOOL

Monterey, California



DENSITY INHOMOGENEITY IN A LASER
CAVITY DUE TO ENERGY RELEASE

by

Allen E. Fuhs

May 1972

Approved for public release; distribution unlimited.

NAVAL POSTGRADUATE SCHOOL
Monterey, California

Rear Admiral A. S. Goodfellow, Jr., USN
Superintendent

Milton U. Clauser
Provost

ABSTRACT:

Density gradients, which refract laser light within the cavity, degrade beam quality. In addition to wall influences and viscous effects which cause density gradients, there is another mechanism. This mechanism, which is due to wakes and compression waves from heat (vibration energy to translation and rotation) addition in a supersonic stream, appears to have been overlooked. The appropriate equation is stated and discussed. A semigraphical solution procedure is outlined. Contours of constant density have been calculated for circular and rectangular cavities. Graphs of the isodensity contours are given.

DENSITY INHOMOGENEITY IN A LASER
CAVITY DUE TO ENERGY RELEASE

by

Dr. Allen E. Fuhs
Professor of Aeronautics
Naval Postgraduate School
Monterey, California 93940

Beam quality is degraded when there are density gradients in the lasing medium. The density gradients refract the light within the cavity. Considerable effort has been devoted by engineers and scientists working with gas dynamic lasers to measure and correct density variations arising from wakes, boundary layers and wall irregularities. Another mechanism exists to cause density gradients, and it appears that this mechanism has been overlooked or ignored.

Heat addition in a supersonic stream causes compression waves which radiate from the energy release region. When the laser radiation is created by stimulated emission (e.g. CO₂ laser), there is also a transfer of vibrational energy to translational and rotational degrees of freedom. This is effectively heat addition and can be treated as a energy-per-unit-volume-and-per-unit-time term in the energy equation.

Using the results from the paper by Tsien and Bielock⁽¹⁾, the following equation can be derived⁽²⁾ for distributed energy sources

$$\frac{\Delta \rho(x,y)}{\rho} = \frac{(\gamma-1)M}{2a^3 \beta \rho} \int_0^S h(x,y) \sin \mu \, dS \quad (1)$$

$$- \frac{\gamma-1}{a^2 U \rho} - \int_{x'} \int_{y'} h(x',y') \delta(y-y') I(x-x') dx' dy'$$

Equation (1) is applicable to a planar geometry; planes normal to the beam are considered. The symbols have the following meaning;

γ	ratio of heat capacity of lasing medium
M	Mach number
a	local speed of sound
β	$(M^2 - 1)^{\frac{1}{2}}$
ρ	density
$h(x', y')$	energy released at (x', y') per unit at volume and unit time
μ	Mach angle equal to arcsin of $1/M$
S	distance along a characteristic
U	local flow velocity
$\delta(y - y')$	delta function
$I(x - x')$	unit function, zero for $x - x' < 0$ and unit for $x - x' > 0$

The first integral yields the density variation along a characteristic. When $h(x, y) = \text{constant}$, the change in density is proportional to the length of characteristic imbedded within the energy release region. When $h(x, y)$ is variable, the element of characteristic length dS is weighted by $h(x, y)$.

The second integral results from the wake of the energy release region. Note that it is opposite in sign to the first integral. The change in density is proportional to the length of streamline imbedded in the energy release zone upstream of the observation point. This is true if $h(x, y)$ is a constant. These facts concerning the proportionality of $\Delta\rho/\rho$ to the length of characteristic and streamline traversing the

energy release region suggest a semigraphical calculation procedure. This is illustrated in Figure 1 for point P with a flow at $M=4$. As illustrated, the right running characteristic R has a length of 7.0, the wake streamline W has a length of 8.5, and the left running characteristic L, a length of 10.2. The length of characteristics must be multiplied by $\sin \mu$. It has been assumed that $h(x,y)$ is a constant within the circle and zero outside the circle. Waves reflected from the walls have been neglected.

It is necessary to evaluate the value of h for typical laser operating conditions. A gas dynamic laser has an efficiency of approximately 1 percent based on the chemical energy which increased the medium temperature from a room value to 1500°K or so. The increase in enthalpy of the gas mixture is 550 BTU/lbm. Consider the flow of gas through a volume with a shape of a cube having 1 ft sides. For typical conditions this gives a flow of 40 lbm/sec through the cube. Assume the cube is the laser cavity and that 1 percent of the energy is removed as radiation. A CO_2 laser has a quantum efficiency of 40%. It is assumed that 1.5% of the energy, which had been frozen in vibration, is transferred to translation and rotation. For these conditions h has a value of 330 BTU/sec ft³. Using appropriate values for M , a , ρ , etc., it is found that

$$\frac{\Delta \rho}{\rho} = 0.023 \text{ per foot of characteristic length}$$

and

$$\frac{\Delta \rho}{\rho} = -0.052 \text{ per foot of streamline in cavity}$$

Using these values of $\Delta \rho/\rho$ per foot of length, the density contours were calculated for both a square and a circular cavity. Figure 2 shows

the results for a circular cavity. Flow is from left to right. At the top of the cavity there is a region of large positive $\Delta\rho/\rho$ due to the fact the wake is small and the left running characteristic within the cavity is long. At the extreme downstream edge of the cavity $\Delta\rho/\rho$ is large negative. For this position there is a long wake within the circle which dominates the compression due to characteristics. At the top of the cavity the 0, +.001, +.002, and +.003 contours have a slope nearly equal to the slope of a left running characteristic. This causes a near loop in the +.003 contour.

In Figure 3 the feature most readily obvious is the slope of the density contours in the upper part of the figure. These might be diagnosed as resulting from waves originating at the upper wall. However, these are due to energy release in a supersonic stream. Along the upper surface the density decreases going downstream. Along the laser centerline the density increases slightly as one moves downstream. In the upper downstream corner there is a strong density gradient.

The analysis of this note has focused on two cases where $h(x,y)$ is constant. For the circular case with a real laser, the value of h is best described as Gaussian. This solution is straightforward but extremely tedious. A computer program seems appropriate for the problem of variable $h(x,y)$.

This note has demonstrated that the energy release can cause $\Delta\rho/\rho$ values somewhat less than viscous flow effects but nonetheless significant. Furthermore, since the $\Delta\rho/\rho$ may be oriented nearly along characteristics, these may be confused with waves originating at a wall.

REFERENCES

1. H. S. Tsien and M. Beilock, "Heat Source in a Uniform Flow," J. Aero Sci, 16, p. 756, 1948.
2. A. E. Fuhs, "Quasi Area Rule for Heat Addition in Transonic and Supersonic Flight Regimes," AF Aero Propulsion Laboratory TR-72-10, WPAFB, Ohio, 1972.

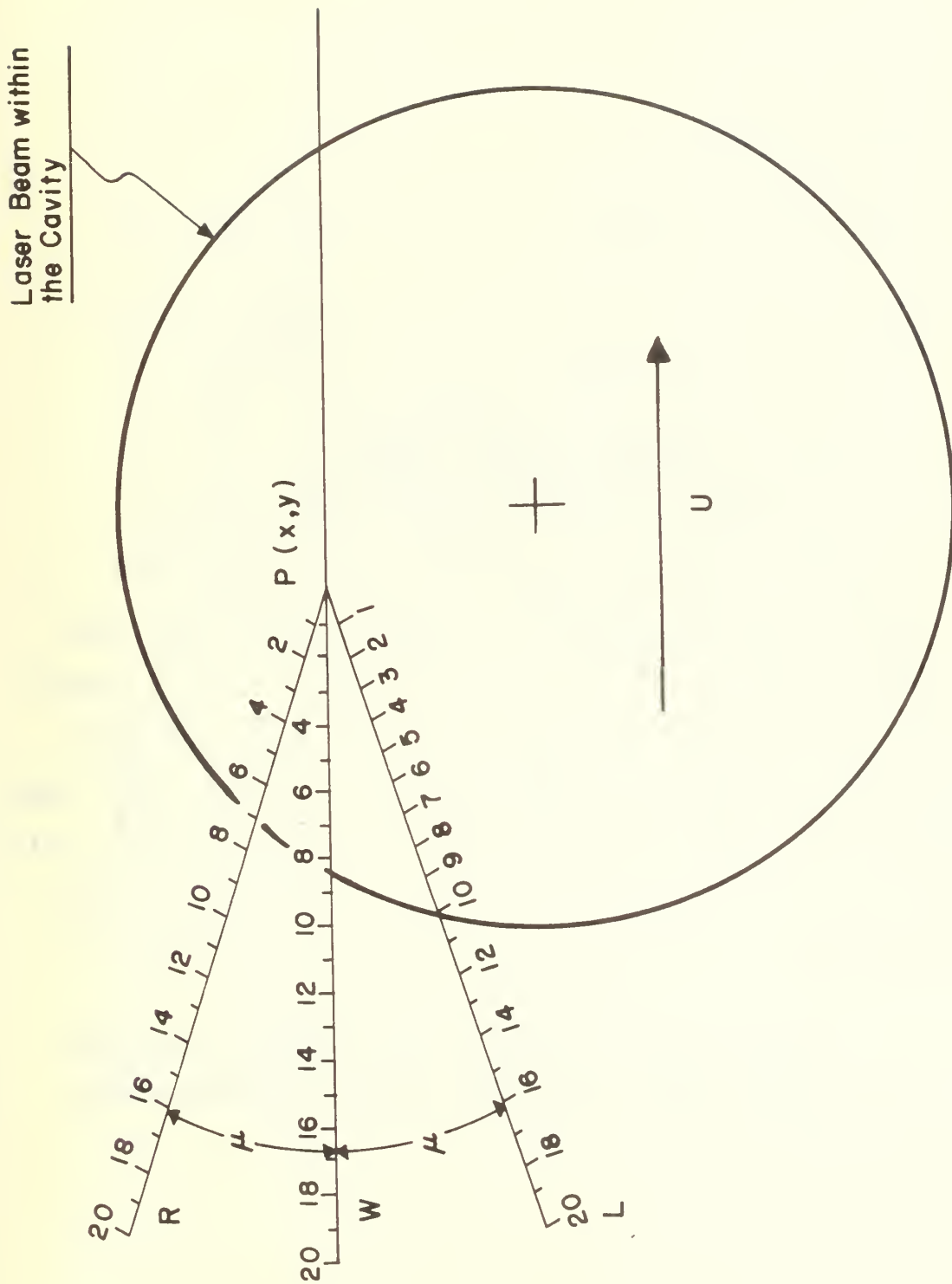


Fig. 1. Illustration of computational procedure.

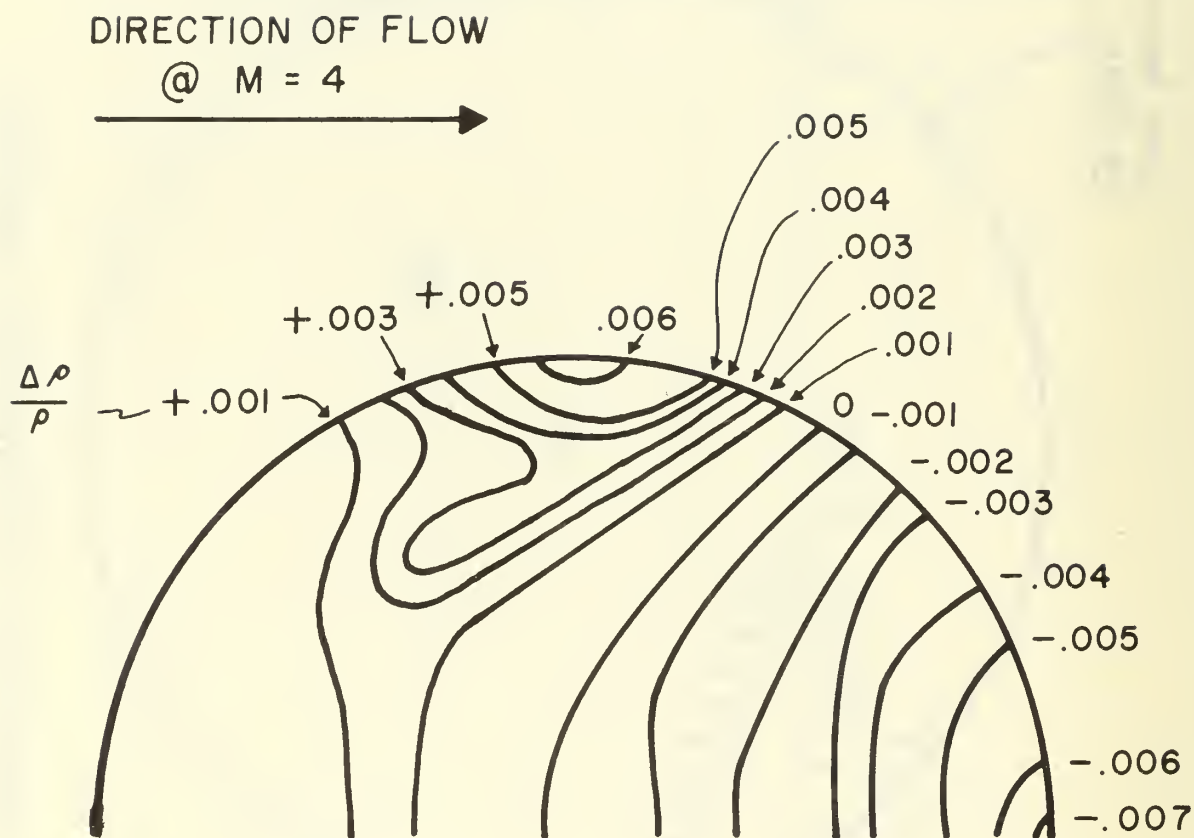


Fig. 2: Fractional density contours due to energy release in a circular laser cavity. Only top half is shown due to symmetry.

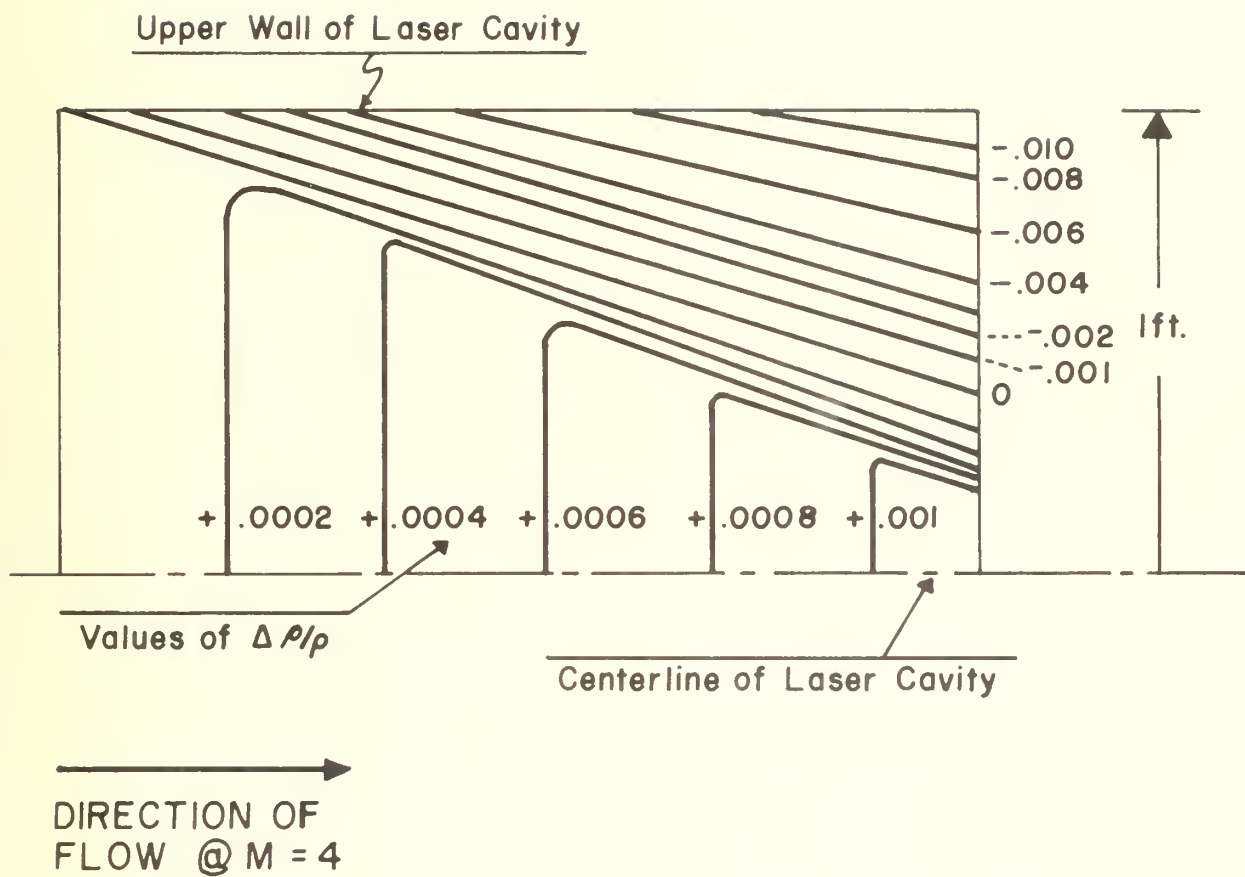


Fig. 3 : Fractional density contours due to energy release in a square cavity.

DISTRIBUTION LIST

	No. Copies
1. Defense Documentation Center Cameron Station Alexandria, Virginia 22314	20
2. Library, Code 0212 Naval Postgraduate School Monterey, California 93940	2
3. Dean of Research Administration Naval Postgraduate School Monterey, California 93940	1
4. Department of Aeronautics Naval Postgraduate School Monterey, California 93940	
Professor R. W. Bell, Chairman	1
Professor R. E. Ball	1
Professor M. Bank	1
Professor J. A. J. Bennett	1
Professor D. J. Collins	1
Professor O. Biblarz	1
Professor A. E. Fuhs	1
Professor T. H. Gawain	1
Professor R. A. Hess	1
Professor G. Hokenson	1
Professor C. H. Kahr	1
Professor D. M. Layton	1
Professor G. H. Lindsey	1
Professor J. A. Miller	1
Professor D. W. Netzer	1
Professor M. F. Platzer	1
Professor H. L. Power	1
Professor M. H. Redlin	1
Professor W. Schlachter	1
Professor L. V. Schmidt	1
Professor R. P. Shreeve	1
Professor M..H. Vavra	1
Professor R. D. Zucker	1
5. Professor G. E. Schacher, Code 61Sq Department of Physics Naval Postgraduate School Monterey, California 93940	1
6. Professor O. Heinz, Code 61Hz Department of Physics Naval Postgraduate School Monterey, California 93940	1

7. Professor N. E. J. Boston, Code 58Bb 1
Department of Oceanography
Naval Postgraduate School
Monterey, California 93940
8. Professor A. Cooper, Code 61Cr 1
Department of Physics
Naval Postgraduate School
Monterey, California 93940
9. Professor Schwirzke, Code 61Sw 1
Department of Physics
Naval Postgraduate School
Monterey, California 93940
10. Professor Ciglio, Code 61 C1 1
Department of Physics
Naval Postgraduate School
Monterey, California 93940
11. Professor Kalmbach, Code 61Kb 1
Department of Physics
Naval Postgraduate School
Monterey, California 93940
12. Professor Neighbours, Code 61Nb 1
Department of Physics
Naval Postgraduate School
Monterey, California 93940
13. Professor Tao, Code 52Tv 1
Department of Electrical Engineering
Naval Postgraduate School
Monterey, California 93940
14. Professor Davidson, Code 51Ds 1
Department of Meteorology
Naval Postgraduate School
Monterey, California 93940
15. Professor Tolles, Code 5417 1
Department of Material Science and Chemistry
Naval Postgraduate School
Monterey, California 93940
16. Professor Powers, Code 52Po 1
Department of Electrical Engineering
Naval Postgraduate School
Monterey, California 93940

17. Gene Crittenden 1
Naval Research Labs
Washington, D. C. 20390
18. Dr. John McCallum 1
Naval Research Labs
Washington, D. C. 20390
19. CAPT James Wilson, USN 1
Naval Ordnance Systems Command
Washington, D. C. 20360
20. Col. Donald Lamberson 1
AFWL
Kirtland AFB, New Mexico
21. Ernest G. Brock 1
North American Rockwell Corp.
Downey, California 90240
22. James R. Carter 1
Naval Weapons Center
China Lake, Calif. 93555
23. Robert J. Collins 1
University of Minnesota
Minneapolis, Minnesota 55400
24. Robert S. Cooper 1
MIT Lincoln Laboratory
Lexington, Massachusetts 02173
25. Jack D. Daugherty 1
AVCO Everett Research Lab
Everett, Massachusetts 02149
26. Anthony J. DeMaria 1
United Aircraft Research Labs
East Hartford, Connecticut 06108
27. George A. Emmons 1
U. S. Army Missile Command
Redstone Arsenal, Alabama 35809
28. William C. Eppers, Jr. 1
Air Force Avionics Lab
Wright-Patterson AFB, Ohio 45433
29. Peter A. Franken 1
The University of Michigan
Ann Arbor, Michigan 48103

30. Harold Jacobs 1
U. S. Army Electronics Command
Fort Monmouth, New Jersey 07703
31. Walter B. Jennings, Jr. 1
U. S. Army Missile Command
Redstone Arsenal, Alabama 35809
32. CAPT Dale A. Holmes 1
Air Force Weapons Laboratory
Kirtland AFB, New Mexico 87117
33. John M. Hood, Jr. 1
Naval Electronics Lab Center
San Diego, California 92100
34. Thomas B. Dowd 1
Executive Secretary
Fifth DoD Conference on Laser Tech
Office of Naval Research
Boston, Massachusetts 02100
35. Frank A. Horrigan 1
Raytheon Research Division
Walham, Massachusetts 02154
36. A. Fenner Milton 1
Naval Research Lab
Washington, D. C. 20390
37. Edwin N. Myers 1
Office of Secretary of Defense
(ODDR&E)
Washington, D. C. 20000
38. Fred W. Quelie, Jr. 1
Office of Naval Research
Boston, Massachusetts 02100
39. Howard R. Schlossberg 1
Air Force Cambridge Research Lab
Bedford, Massachusetts 01730
40. William C. Schoonover 1
Air Force Avionics Lab
Wright-Patterson AFB, Ohio 45433
41. Walter R. Sooy 1
Naval Research Laboratory
Washington, D. C. 20390

42. C. Martin Stickley 1
Advanced Research Projects Agency
Arlington, Virginia 22200
43. Malcolm L. Stitch 1
Union Carbide Corporation
Korad Division
Santa Monica, California 90406
44. Robert B. Watson 1
Office of Chief of Research and Development
Department of the Army
Washington, D. C. 20000
45. Eric J. Woodbury 1
Hughes Aircraft Company
Culver City, California 90230
46. George J. Zissis 1
The University of Michigan
Ann Arbor, Michigan 48103

DOCUMENT CONTROL DATA - R & D

(Security classification of title, body of abstract and indexing annotation must be entered when the overall report is classified)

1. ORIGINATING ACTIVITY (Corporate author) Naval Postgraduate School Monterey, California 93940		2a. REPORT SECURITY CLASSIFICATION Unclassified	
		2b. GROUP	
3. REPORT TITLE Density Inhomogeneity in a Laser Cavity Due to Energy Release			
4. DESCRIPTIVE NOTES (Type of report and, inclusive dates)			
5. AUTHOR(S) (First name, middle initial, last name) Allen E. Fuhs			
6. REPORT DATE May 1972		7a. TOTAL NO. OF PAGES 16	7b. NO. OF REFS 2
8a. CONTRACT OR GRANT NO.		9a. ORIGINATOR'S REPORT NUMBER(S) NPS-57Fu72051A	
b. PROJECT NO.			
c.		9b. OTHER REPORT NO(S) (Any other numbers that may be assigned this report)	
d.			
10. DISTRIBUTION STATEMENT This document has been approved for public release and sale; its distribution is unlimited.			
11. SUPPLEMENTARY NOTES		12. SPONSORING MILITARY ACTIVITY Naval Postgraduate School Monterey, California 93940	
13. ABSTRACT Density gradients, which refract laser light within the cavity, degrade beam quality. In addition to wall influences and viscous effects which cause density gradients, there is another mechanism. This mechanism, which is due to wakes and compression waves from heat (vibration energy to translation and rotation) addition in a supersonic stream, appears to have been overlooked. This appropriate equation is stated and discussed. A semigraphical solution procedure is outlined. Contours of constant density have been calculated for circular and rectangular cavities. Graphs of the isodensity contours are given.			

14 KEY WORDS	LINK A		LINK B		LINK C	
	ROLE	WT	ROLE	WT	ROLE	WT
Gas dynamic lasers Electrical lasers Beam quality Supersonic heat addition						

DUDLEY KNOX LIBRARY - RESEARCH REPORTS



5 6853 01058143 2

